



Knot Invariants from Zero-Dimensional QFT

Abstract. For the purpose of today, an “I-Type Knot Invariant” is a knot invariant computed from a knot diagram by integrating the exponential of a *perturbed Gaussian* Lagrangian which is a sum over the features of that diagram (crossings, edges, faces) of locally defined quantities, over a product of finite dimensional spaces associated to those same features.

joint with
R. van der Veen

Q. Are there any such things? **A.** Yes.

Q. Are they any good? **A.** They are the strongest we know per CPU cycle, and are excellent in other ways too.

Q. Didn’t Witten do that back in 1988 with path integrals?

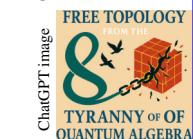
A. No. His constructions are infinite dimensional and far from rigorous.

Q. But integrals belong in analysis!

A. Ours only use squeaky-clean algebra.

Dreams. Given a knot K with a Seifert surface Σ , we dream that there is a 0D Lagrangian $L_\Sigma : 6H_1(\Sigma; \mathbb{R}) \rightarrow \mathbb{R}$ whose coefficients are (low degree) finite type invariants of graphs representing multiple homology classes, such that

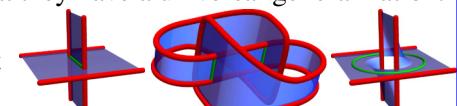
$$Z = \oint_{6H_1(\Sigma; \mathbb{R})} \exp(L_\Sigma)$$



is the invariant θ presented below. We dream the formulas are simple and natural, and that they have a universal generalization.

Sweet. Even sweeter:

May say something about ribbon knots!



The $sl_2^{/\epsilon^2}$ Example. With T an indeterminate and with $\epsilon^2 = 0$:

$$\text{Diagram showing a knot with various components labeled with } L(X_{ij}^+), L(C_i^{-1}), \text{ and } R^2_{p_i x_j} \text{ terms.} \rightarrow Z = \oint_{\mathbb{R}^{14}_{p_i x_i}} \mathcal{L}(X_{15}^+) \mathcal{L}(X_{62}^+) \mathcal{L}(X_{37}^+) \mathcal{L}(C_4^{-1})$$

where $\mathcal{L}(X_{ij}^s) = T^{s/2} \mathbb{E}^{L(X_{ij}^s)}$ and $\mathcal{L}(C_i^\varphi) = T^{\varphi/2} \mathbb{E}^{L(C_i^\varphi)}$, and

$$\begin{aligned} L(X_{ij}^s) &= x_i(p_{i+1} - p_i) + x_j(p_{j+1} - p_j) \\ &\quad + (T^s - 1)x_i(p_{i+1} - p_{j+1}) \\ &\quad + \frac{\epsilon s}{2} \left(x_i(p_i - p_j) \left(\begin{matrix} (T^s - 1)x_i p_j \\ + 2(1 - x_j p_j) \end{matrix} \right) - 1 \right) \\ L(C_i^\varphi) &= x_i(p_{i+1} - p_i) + \epsilon \varphi (1/2 - x_i p_i) \end{aligned}$$

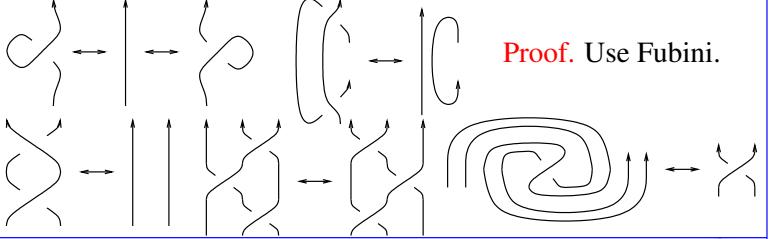
So $Z = T \oint \mathbb{E}^{L(\otimes)} dp_1 \dots dp_7 dx_1 \dots dx_7$, where $L(\otimes)$

$$\sum_{i=1}^7 x_i(p_{i+1} - p_i) + (T-1)(x_1(p_2 - p_6) + x_6(p_7 - p_3) + x_3(p_4 - p_8)) + \frac{\epsilon}{2} \left(\begin{array}{l} x_1(p_1 - p_5)((T-1)x_1 p_5 + 2(1 - x_5 p_5)) - 1 \\ + x_6(p_6 - p_2)((T-1)x_6 p_2 + 2(1 - x_2 p_2)) - 1 \\ + x_3(p_3 - p_7)((T-1)x_3 p_7 + 2(1 - x_7 p_7)) - 1 \\ + 2x_4 p_4 - 1 \end{array} \right)$$

and so $Z = (T - 1 + T^{-1})^{-1} \exp\left(\epsilon \cdot \frac{(T-2+T^{-1})(T+T^{-1})}{(T-1+T^{-1})^2}\right) = \Delta^{-1} \exp\left(\epsilon \cdot \frac{(T-2+T^{-1})\rho_1}{\Delta^2}\right)$. Here Δ is the Alexander polynomial and ρ_1 is the Rozansky-Overbay polynomial [R1, R2, R3, Ov, BV1, BV2].



Theorem. Z is a knot invariant.



Proof. Use Fubini.



(Alternative) Gaussian Integration.

Goal. Compute $\int_{\mathbb{R}^n} dx \exp\left(-\frac{1}{2} a^{ij} x_i x_j + V(x)\right)$.

Solution. Set $\mathcal{Z}_\lambda(x) := \lambda^{-n/2} \int_{\mathbb{R}^n} dy \exp\left(-\frac{1}{2\lambda} a^{ij} y_i y_j + V(x+y)\right)$.

Then $\mathcal{Z}_1(0)$ is what we want, $\mathcal{Z}_0(x) = (\det A)^{-1/2} \exp V(x)$, and with g_{ij} the inverse matrix of a^{ij} and noting that under the dy

integral $\partial_y = 0$,

$$\begin{aligned} &= \frac{\lambda^{-n/2}}{2} \int_{\mathbb{R}^n} dy g_{ij}(\partial_{x_i} - \partial_{y_i})(\partial_{x_j} - \partial_{y_j}) \exp\left(-\frac{1}{2\lambda} a^{ij} y_i y_j + V(x+y)\right) \\ &= \frac{\lambda^{-n/2}}{2\lambda^2} \int_{\mathbb{R}^n} dy (g_{ij} a^{ii'} a^{jj'} y_i y_j - \lambda g_{ij} a^{ji'}) \exp\left(-\frac{1}{2\lambda} a^{ij} y_i y_j + V(x+y)\right) \\ &= \frac{\lambda^{-n/2}}{2\lambda^2} \int_{\mathbb{R}^n} dy (a^{ij} y_i y_j - \lambda n) \exp\left(-\frac{1}{2\lambda} a^{ij} y_i y_j + V(x+y)\right) \\ &= \partial_\lambda \mathcal{Z}_\lambda(x). \end{aligned}$$

Hence

$$(*) \quad \partial_\lambda \mathcal{Z}_\lambda(x) = \frac{1}{2} g_{ij} \partial_{x_i} \partial_{x_j} \mathcal{Z}_\lambda(x),$$

and therefore $\mathcal{Z}_\lambda(x) = (\det A)^{-1/2} \exp\left(\frac{\lambda}{2} g_{ij} \partial_{x_i} \partial_{x_j}\right) \exp V(x)$.

We’ve just witnessed the birth of “Feynman Diagrams”.

Even better. With $Z_\lambda := \log(\sqrt{\det A} \mathcal{Z}_\lambda)$, by a simple substitution into (*), we get the “Synthesis Equation”:

$Z_0 = V, \quad \partial_\lambda Z_\lambda = \frac{1}{2} \sum_{i,j=1}^n g_{ij} (\partial_{x_i} Z_\lambda + (\partial_{x_i} Z_\lambda)(\partial_{x_j} Z_\lambda)) =: F(Z_\lambda)$, an ODE (in λ) whose solution is pure algebra.



Picard Iteration (used to prove the existence and uniqueness of solutions of ODEs). To solve $\partial_\lambda f_\lambda = F(f_\lambda)$ with a given f_0 , start with f_0 , iterate $f \mapsto f_0 + \int_0^\lambda F(f_\lambda) d\lambda$, and seek a fixed point. In our cases, it is always reached after finitely many iterations!



Definition. \oint : The result of this process, ignoring the convergence of the actual integral.

Strong. The pair (Δ, ρ_1) attains 270,316 distinct values on the 313,230 prime knots with up to 15 crossings (a deficit of 42,914), whereas the pair $(H = \text{HOMFLYPT polynomial}, Kh = \text{Khovanov Homology})$ attains only 242,985 distinct values on the same knots (a deficit of 70,245). The pair (Δ, θ) , discussed later, has a deficit of only 6,758, and the triple (Δ, θ, ρ_2) , of only 6,341.

Yet better than (H, Kh) and other Reshetikhin-Turaev-Witten invariants and knot homologies, Δ, ρ_1, ρ_2 and θ can be computed in **polynomial time** (and hence, even for very large knots).

So ugly as the formulas may be (and θ ’s formulas are uglier), these invariants are possibly **the best we have!**

Acknowledgement. This work was supported by NSERC grants RGPIN-2018-04350 and RGPIN-2025-06718 and by the Chu Family Foundation (NYC).

Implementation (see `IType.nb` of $\omega\epsilon\beta/\text{ap}$).

⊕ Once[`<< KnotTheory``; `<< Rot.m`];

□ Loading `KnotTheory`` version
of October 29, 2024, 10:29:52.1301.

Read more at <http://katlas.org/wiki/KnotTheory>.

□ Loading `Rot.m` from <http://drorbn.net/AP/Talks/Bonn-2505>
to compute rotation numbers.

⊕ `CF[ω_. ε_. E] := CF[ω] × CF /@ ε;`
`CF[ε_. List] := CF /@ ε;`
`CF[ε_] := Module[{vs, ps, c},`
 `vs = Cases[ε, (x | p | ε | π | g) __, ∞] ∪ {e};`
 `Total[CoefficientRules[Expand[ε], vs] /.`
 `(ps_ → c_) ↦ Factor[c] (Times @@ vs^ps)]];`

Integration using Picard iteration. The core is in yellow and hacks are in pink.

⊕ `E /: E[A_] × E[B_] := E[A + B];`

⊕ `$π = Identity; (* The Wisdom Projection *)`

⊕ `Unprotect[Integrate];`

□ $\int \omega_ \cdot E[L_] d(vs_List) :=$

`Module[{n, L0, Q, Δ, G, Z0, Z, λ, DZ, DDZ,`
 `FZ, a, b},`
`n = Length@vs; L0 = L /. e → 0;`
`Q = Table[(-Δ vs[[a]], vs[[b]] L0) /. Thread[vs → 0] /.`
 `(p | x) __ → 0, {a, n}, {b, n}];`
`If[(Δ = Det[Q]) == 0, Return@"Degenerate Q!"];`
`Z = Z0 = CF[$π(L + vs.Q.vs / 2); G = Inverse[Q];`
`FixedPoint[(DZ = Table[∂v Z, {v, vs}];`
 `DDZ = Table[∂u DZ, {u, vs}];`
 `FZ = Sum[G[[a, b]] (DDZ[[a, b]] + DZ[[a]] × DZ[[b]]),`
 `{a, n}, {b, n}] / 2;`
 `Z = CF[Z0 + Integrate[$π FZ dλ, {λ, 0, π}], &, Z];`
 `PowerExpand@Factor[ω Δ⁻¹/²] ×`
 `E[CF[Z /. λ → 1 /. Thread[vs → 0]]]];`
`Protect[Integrate];`

⊕ $\int E[-μ x^2 / 2 + i ε x] d\{x\}$

□ $\frac{E\left[-\frac{\xi^2}{2\mu}\right]}{\sqrt{\mu}}$

⊕ $FoF G = \int E[-μ (x - a)^2 / 2 + i ε x] d\{x\}$

□ $\frac{E\left[\frac{i(2a\mu + iε)\xi}{2\mu}\right]}{\sqrt{\mu}}$

⊕ $\int FoF G E[-i ε x] d\{ε\}$

□ $E\left[-\frac{1}{2} (a - x)^2 \mu\right]$

So we've tested and nearly proven the Fourier inversion formula!

$$\odot L = -\frac{1}{2} \{x_1, x_2\} \cdot \begin{pmatrix} a & b \\ b & c \end{pmatrix} \cdot \{x_1, x_2\} + \{\xi_1, \xi_2\} \cdot \{x_1, x_2\};$$

$$\odot Z_{12} = \int E[L] d\{x_1, x_2\}$$

$$\square \frac{E\left[\frac{c \xi_1^2}{2 (-b^2 + a c)} + \frac{b \xi_1 \xi_2}{b^2 - a c} + \frac{a \xi_2^2}{2 (-b^2 + a c)}\right]}{\sqrt{-b^2 + a c}}$$

$$\odot \{Z_1 = \int E[L] d\{x_1\}, Z_{12} = \int Z_1 d\{x_2\}\}$$

$$\square \left\{ \frac{E\left[-\frac{(-b^2 + a c) x_2^2}{2 a} - \frac{b x_2 \xi_1}{a} + \frac{\xi_1^2}{2 a} + x_2 \xi_2\right]}{\sqrt{a}}, \text{True} \right\}$$

$$\odot \$π = \text{Normal}[\# + 0[e]^{13}] \&; \int E[-φ^2/2 + e φ^3/6] d\{φ\}$$

$$\square E\left[\frac{5 e^2}{24} + \frac{5 e^4}{16} + \frac{1105 e^6}{1152} + \frac{565 e^8}{128} + \frac{82825 e^{10}}{3072} + \frac{19675 e^{12}}{96}\right]$$



Guido Fubini

From <https://oeis.org/A226260>:

0 1 3 6 2 7
 OEIS
 THE ON-LINE ENCYCLOPEDIA
 OF INTEGER SEQUENCES®
 23 20 12
 10 22 11 21

founded in 1964 by N. J. A. Sloane

 (Greetings from The On-Line Encyclopedia of Integer Sequences!)

A226260 Numerators of mass formula for connected vacuum graphs on $2n$ nodes for a ϕ^3 field theory.
 1, 5, 5, 1185, 565, 82825, 19675, 1282031525, 80727925, 168348621875, 13209845125,
 2239646759308375, 19739117698375, 6326791709083309375, 32468078556378125, 38362676768845045751875,
 281365778405032973125, 2824650747089425586152484375, 77663215703416712734375 ([list](#); [graph](#); [refs](#); [listen](#); [history](#); [text](#); [internal format](#))

The Right-Handed Trefoil.

⊕ `K = Mirror@Knot[3, 1]; Features[K]`

□ `Features[7, C4[-1] X1,5[1] X3,7[1] X6,2[1]]`

⊕ $L[X_{i,j}[s_]] := T^{s/2} E\left[x_i(p_{i+1} - p_i) + x_j(p_{j+1} - p_j) + (T^s - 1)x_i(p_{i+1} - p_{j+1}) + (e s / 2) \times (x_i(p_i - p_j) ((T^s - 1)x_i p_j + 2(1 - x_j p_j)) - 1)\right]$
 $L[C_{i, [φ]}] := T^{φ/2} E\left[x_i(p_{i+1} - p_i) + e φ \left(\frac{1}{2} - x_i p_i\right)\right]$

`L[K_] := CF[L /@ Features[K][2]]`

`vs[K_] :=`

`Join @@ Table[{pi, xi}, {i, Features[K][1]}]`

⊕ `{vs[K], L[K]}`

□ $\{p_1, x_1, p_2, x_2, p_3, x_3, p_4, x_4, p_5, x_5, p_6, x_6, p_7, x_7\},$
 $T E\left[-2e - p_1 x_1 + e p_1 x_1 + T p_2 x_1 - e p_5 x_1 + (1 - T) p_6 x_1 + \frac{1}{2} (-1 + T) e p_1 p_5 x_1^2 + \frac{1}{2} (1 - T) e p_5^2 x_1^2 - p_2 x_2 + p_3 x_2 - p_3 x_3 + e p_3 x_3 + T p_4 x_3 - e p_7 x_3 + (1 - T) p_8 x_3 + \frac{1}{2} (-1 + T) e p_3 p_7 x_3^2 + \frac{1}{2} (1 - T) e p_7^2 x_3^2 - p_4 x_4 + e p_4 x_4 + p_5 x_4 - p_5 x_5 + p_6 x_5 - e p_1 p_5 x_1 x_5 + e p_5^2 x_1 x_5 - e p_2 x_6 + (1 - T) p_3 x_6 - p_6 x_6 + e p_6 x_6 + T p_7 x_6 + e p_2^2 x_2 x_6 - e p_2 p_6 x_2 x_6 + \frac{1}{2} (1 - T) e p_2^2 x_6^2 + \frac{1}{2} (-1 + T) e p_2 p_6 x_6^2 - p_7 x_7 + p_8 x_7 - e p_3 p_7 x_3 x_7 + e p_7^2 x_3 x_7\}$



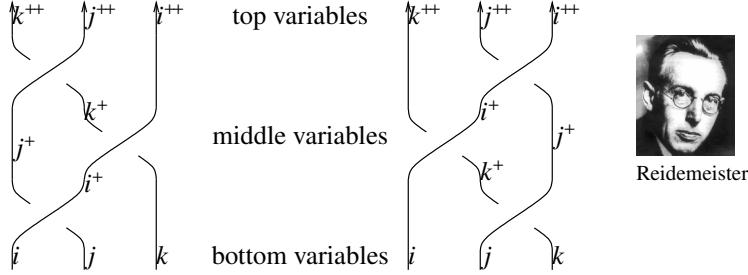
Joseph Fourier

$$\textcircled{S} \pi = \text{Normal}[\# + 0[\epsilon]^2] \& \int \mathcal{L}[K] d\pi[K]$$

$$\square - \frac{\frac{1}{2} T \mathbb{E} \left[- \frac{(-1+T)^2 (1+T^2) \epsilon}{(1-T+T^2)^2} \right]}{1-T+T^2}$$

A faster program to compute ρ_1 , and more stories about it, are at [BV2].

Invariance Under Reidemeister 3.



$\textcircled{S} \text{lhs}$

$$\square T^{3/2} \mathbb{E} \left[- \frac{3 \epsilon}{2} + i T^2 p_{2+i} \pi_i - i (-1+T) T p_{2+j} \pi_i + i T^2 p_{2+j} \pi_i - i (-1+T) p_{2+k} \pi_i + \frac{1}{2} T p_{2+k} \pi_i - \frac{1}{2} (-1+T) T^3 p_{2+i} p_{2+j} \pi_i^2 + \frac{1}{2} (-1+T) T^3 p_{2+j} \pi_i^2 - \frac{1}{2} (-1+T) T^2 p_{2+i} p_{2+k} \pi_i^2 + \frac{1}{2} (-1+T)^2 T p_{2+j} p_{2+k} \pi_i^2 + \frac{1}{2} (-1+T) T p_{2+k} \pi_i^2 + i T p_{2+j} \pi_j - i T p_{2+j} \pi_j - i (-1+T) p_{2+k} \pi_j + i (-1+2T) p_{2+k} \pi_j + T^3 p_{2+i} p_{2+j} \pi_i \pi_j - T^3 p_{2+j} \pi_i \pi_j - (-1+T) T^2 p_{2+i} p_{2+k} \pi_i \pi_j + (-1+T)^2 T p_{2+j} p_{2+k} \pi_i \pi_j + (-1+T) T p_{2+k} \pi_i \pi_j - \frac{1}{2} (-1+T) T p_{2+j} p_{2+k} \pi_j^2 + \frac{1}{2} (-1+T) T p_{2+k} \pi_j^2 + i p_{2+k} \pi_k - 2 i p_{2+k} \pi_k + T^2 p_{2+i} p_{2+k} \pi_i \pi_k - (-1+T) T p_{2+j} p_{2+k} \pi_i \pi_k - T p_{2+k} \pi_i \pi_k + T p_{2+j} p_{2+k} \pi_j \pi_k - T p_{2+k} \pi_j \pi_k \right]$$

Invariance under the other Reidemeister moves is proven in a similar way. See IType.nb at [ωεβ/ap](#).

$$\textcircled{S} \text{lhs} = \int (\mathcal{L} /@ (X_{i,j}[1] X_{i+1,k}[1] X_{j+1,k+1}[1]))$$

$$d\{p_{i+1}, p_{j+1}, p_{k+1}, x_{i+1}, x_{j+1}, x_{k+1}\};$$

$$\text{rhs} = \int (\mathcal{L} /@ (X_{j,k}[1] X_{i,k+1}[1] X_{i+1,j+1}[1]))$$

$$d\{x_{i+1}, p_{i+1}, p_{j+1}, p_{k+1}, x_{j+1}, x_{k+1}\};$$

$$\text{lhs} === \text{rhs}$$

\square False

Invariance Under Reidemeister 3, Take 2.

$$\textcircled{S} \text{lhs} = \int (\mathcal{L} /@ (X_{i,j}[1] X_{i+1,k}[1] X_{j+1,k+1}[1]))$$

$$d\{x_i, x_j, x_k, p_{i+1}, p_{j+1}, p_{k+1}, x_{i+1}, x_{j+1}, x_{k+1}\};$$

$$\text{rhs} = \int (\mathcal{L} /@ (X_{j,k}[1] X_{i,k+1}[1] X_{i+1,j+1}[1]))$$

$$d\{x_i, x_j, x_k, x_{i+1}, p_{i+1}, p_{j+1}, p_{k+1}, x_{j+1}, x_{k+1}\};$$

$$\text{lhs} === \text{rhs}$$

\square True

$\textcircled{S} \text{lhs}$

\square Degenerate Q!

$\textcircled{S} \epsilon^2 r_2[1, i, j]$

$$\square \frac{1}{12} \epsilon^2 (-6 p_i x_i + 6 p_j x_i - 3 (-1+3T) p_i p_j x_i^2 + 3 (-1+3T) p_j^2 x_i^2 + 4 (-1+T) p_i^2 p_j x_i^3 - 2 (-1+T) (5+T) p_i p_j^2 x_i^3 + 2 (-1+T) (3+T) p_j^3 x_i^3 + 18 p_i p_j x_i x_j - 18 p_j^2 x_i x_j - 6 p_i^2 p_j x_i^2 x_j + 6 (2+T) p_i p_j^2 x_i^2 x_j - 6 (1+T) p_j^3 x_i^2 x_j - 6 p_i p_j^2 x_i x_j^2 + 6 p_j^3 x_i x_j^2)$$

$\textcircled{S} \epsilon^2 r_2[-1, i, j]$

$$\square \frac{1}{12 T^2} \epsilon^2 (-6 T^2 p_i x_i + 6 T^2 p_j x_i + 3 (-3+T) T p_i p_j x_i^2 - 3 (-3+T) T p_j^2 x_i^2 - 4 (-1+T) T p_i^2 p_j x_i^3 + 2 (-1+T) (1+5T) p_i p_j^2 x_i^3 - 2 (-1+T) (1+3T) p_j^3 x_i^3 + 18 T^2 p_i p_j x_i x_j - 18 T^2 p_j^2 x_i x_j - 6 T^2 p_i^2 p_j x_i^2 x_j + 6 T (1+2T) p_i p_j^2 x_i^2 x_j - 6 T (1+T) p_j^3 x_i^2 x_j - 6 T^2 p_i p_j^2 x_i x_j^2 + 6 T^2 p_j^3 x_i x_j^2)$$

$\textcircled{S} \epsilon^2 \gamma_2[\phi, i]$

$$\square -\frac{1}{2} \epsilon^2 \phi^2 p_i x_i$$

Even more! • The sl_2 formulas mod ϵ^4 are in the last page of the handout of [BN3].

- Using [GPV] we can show that every finite type invariant is I-Type.
- Probably, $\langle \text{Reshetikhin-Turaev} \rangle \subset \langle \text{I-Type} \rangle$ efficiently.
- Possibly, $\langle \text{Rozansky Polynomials} \rangle \subset \langle \text{I-Type} \rangle$ efficiently.
- Knot signatures are I-Type, at least mod 8.
- We already have some work on sl_3 , and it leads to the strongest genuinely-computable knot invariant presently known.

The $s_3^{e^2}$ Example [BV3]. Here we have two formal variables $T_1 \odot F_1[\{s_, i_, j\}__]$:= CF[
and T_2 , we set $T_3 := T_1 T_2$, we integrate over 6 variables for each
edge: $p_{1i}, p_{2i}, p_{3i}, x_{1i}, x_{2i}$, and x_{3i} .

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 $\odot T_3 = T_1 T_2; \quad i\_^+ := i + 1;$ 
$π =
  (CF@Normal[# + 0[e]^2] /.
    {πis\_ → B^-1 πis, xis\_ → B^-1 xis, pis\_ → B pis} /.
    e B^b /; b < 0 → 0 /. B → 1) &;

```

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 $\odot vs_{i\_} := Sequence[p_{1,i}, p_{2,i}, p_{3,i}, x_{1,i}, x_{2,i}, x_{3,i}];$ 
F[is_] := E[Sum[πv,i p_{v,i}, {i, {is}}, {v, 3}]];
L[K_] := CF[L /@ Features[K][2]];
vs[K_] :=
  Union @@ Table[{vs[i], {i, Features[K][1]}}
The Lagrangian.

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 $\odot L[X_{i,j}[s\_\_]] := T_3^s E[CF@Plus[$ 
 $\sum_{i=1}^3 (x_{vi} (p_{vi^+} - p_{vi}) + x_{vj} (p_{vj^+} - p_{vj}) + (T_v^s - 1) x_{vi} (p_{vi^+} - p_{vj^+})),$ 
 $(T_1^s - 1) p_{3j} x_{1i} (T_2^s x_{2i} - x_{2j}),$ 
 $e s (T_3^s - 1) p_{1j} (p_{2i} - p_{2j}) x_{3i} / (T_2^s - 1),$ 
 $e s (1/2 + T_2^s p_{1i} p_{2j} x_{1i} x_{2i} - p_{1i} p_{2j} x_{1i} x_{2j} - p_{3i} x_{3i} -$ 
 $(T_2^s - 1) p_{2j} p_{3i} x_{2i} x_{3i} + (T_3^s - 1) p_{2j} p_{3j} x_{2i} x_{3i} +$ 
 $2 p_{2j} p_{3i} x_{2j} x_{3i} + p_{1i} p_{3j} x_{1i} x_{3j} - p_{2i} p_{3j} x_{2i} x_{3j} -$ 
 $T_2^s p_{2j} p_{3j} x_{2i} x_{3j} +$ 
 $((T_1^s - 1) p_{1j} x_{1i} (T_2^s p_{2j} x_{2i} - T_2^s p_{2j} x_{2j} -$ 
 $(T_2^s + 1) (T_3^s - 1) p_{3j} x_{3i} + T_2^s p_{3j} x_{3j}) +$ 
 $(T_3^s - 1) p_{3j} x_{3i} (1 - T_2^s p_{1i} x_{1i} + p_{2i} x_{2j} + (T_2^s - 2) p_{2j} x_{2j})) /$ 
 $(T_2^s - 1)))]$ 

```

```

 $\odot L[C_{i\_}[\varphi\_\_]] := T_3^\varphi E[\sum_{i=1}^3 x_{vi} (p_{vi^+} - p_{vi}) + e \varphi (p_{3i} x_{3i} - 1/2)]$ 
Reidemeister 3.

```

```

 $\odot Short[$ 
  lhs = ∫ F[i, j, k] × L /@ (X[i,j][1] X[i+,k][1] X[i+,k+][1])
  d{vs[i], vs[j], vs[k], vs[i+], vs[j+], vs[k+]}

```

```

 $\square T_1^3 T_2^3$ 
E[ $\frac{3 e}{2} + T_1^2 p_{1,2+i} \pi_{1,i} - (-1 + T_1) T_1 p_{1,2+j} \pi_{1,i} + <<150>>$ ]

```

```

 $\odot rhs = \int F[i, j, k] \times L /@ (X[j,k][1] X[i,k+][1] X[i+,j+][1])$ 
d{vs[i], vs[j], vs[k], vs[i+], vs[j+], vs[k+]};
lhs == rhs

```

True

The Trefoil.

```

 $\odot K = Knot[3, 1]; \quad \int L[K] d vs[K]$ 
 $\square - \left( \frac{1}{2} T_1^2 T_2^2 \right.$ 
 $E \left[ - \left( \left( 1 - T_1 + T_1^2 - T_2 - T_1^3 T_2 + T_2^2 + T_1^4 T_2^2 - T_1 T_2^3 - \right. \right. \right.$ 
 $T_1^4 T_2^3 + T_1^2 T_2^4 - T_1^3 T_2^4 + T_1^4 T_2^4 \left. \left. \left. \right) \right) / \left( \left( 1 - T_1 + T_1^2 \right) \left( 1 - T_2 + T_2^2 \right) \left( 1 - T_1 T_2 + T_1^2 T_2^2 \right) \right) \right) /$ 
 $\left( \left( 1 - T_1 + T_1^2 \right) \left( 1 - T_2 + T_2^2 \right) \left( 1 - T_1 T_2 + T_1^2 T_2^2 \right) \right)$ 

```



$$\begin{aligned}
 & s (1/2 - g_{3ii} + T_2^s g_{1ii} g_{2ji} - g_{1ii} g_{2jj} - (T_2^s - 1) g_{2ji} g_{3ii} + \\
 & 2 g_{2jj} g_{3ii} - (1 - T_3^s) g_{2ji} g_{3ji} - g_{2ii} g_{3jj} - T_2^s g_{2ji} g_{3jj} + \\
 & g_{1ii} g_{3jj} + \\
 & ((T_1^s - 1) g_{1ji} (T_2^s g_{2ji} - T_2^s g_{2jj} + T_2^s g_{3jj}) + \\
 & (T_3^s - 1) g_{3ji} (1 - T_2^s g_{1ii} - (T_1^s - 1) (T_2^s + 1) g_{1ji} + \\
 & (T_2^s - 2) g_{2jj} + g_{2ij})) / (T_2^s - 1))
 \end{aligned}$$

```

 $\odot F_2[\{s0\_, i0\_, j0\_\}, \{s1\_, i1\_, j1\_\}] :=$ 
CF[s1 (T1^s0 - 1) (T2^s1 - 1)^-1 (T3^s1 - 1) g1,j1,i0 g3,j0,i1
((T2^s0 g2,i1,i0 - g2,i1,j0) - (T2^s0 g2,j1,i0 - g2,j1,j0)))]

```

```

 $\odot F_3[\varphi\_, k\_] = \varphi g_{3kk} - \varphi / 2;$ 

```

We call the invariant computed θ :

```

 $\odot \Theta[K\_] := \Theta[K] = Module[\{X, \varphi, n, A, \Delta, G, ev, \theta\},$ 
  {X, \varphi} = Rot[K];
  n = Length[X];
  A = IdentityMatrix[2 n + 1];
  Cases[X, \{s\_, i\_, j\_\} \rightarrow
    (A[[i, j], {i + 1, j + 1}] += {{-T^s, T^s - 1}, {0, -1}})];
  Δ = T^{(-Total[\varphi] - Total[X[[All, 1]]])/2} Det[A];
  G = Inverse[A];
  ev[\&_] := Factor[\& /. g_{v\_, \alpha\_, \beta\_\_} \rightarrow (G[[\alpha, \beta]] /. T \rightarrow T_v)];
  θ = ev[Sum_{k=1}^n F_1[X[[k]]]];
  θ += ev[Sum_{k1=1}^n Sum_{k2=1}^n F_2[X[[k1]], X[[k2]]]];
  θ += ev[Sum_{k=1}^{2 n} F_3[\varphi[[k]], k]];
  Factor@{\Delta, (\Delta /. T \rightarrow T_1) (\Delta /. T \rightarrow T_2) (\Delta /. T \rightarrow T_3) \θ}]
];

```

Some Knots.

```

 $\odot Expand[\Theta[Knot[3, 1]]]$ 

```

$$\begin{aligned}
 & \square \left\{ -1 + \frac{1}{T} + T, -\frac{1}{T_1^2} - T_1^2 - \frac{1}{T_2^2} - \frac{1}{T_1^2 T_2^2} + \frac{1}{T_1 T_2^2} + \right. \\
 & \left. \frac{1}{T_1^2 T_2} + \frac{T_1}{T_2} + \frac{T_2}{T_1} + T_1^2 T_2 - T_2^2 + T_1 T_2^2 - T_1^2 T_2^2 \right\}
 \end{aligned}$$

```

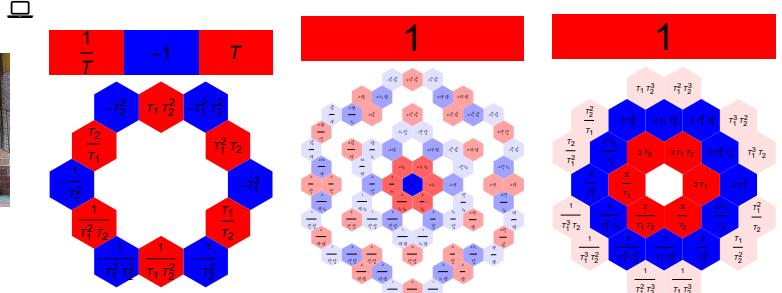
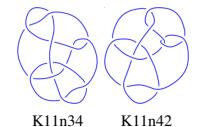
 $\odot (* PolyPlot suppressed *)$ 

```

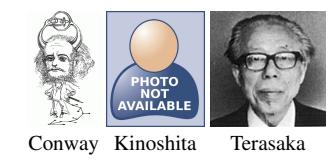
```

 $\odot GraphicsRow[PolyPlot[\Theta[Knot[\#]],$ 
  Labeled → True] &
  /@ {"3_1", "K11n34", "K11n42"}]

```



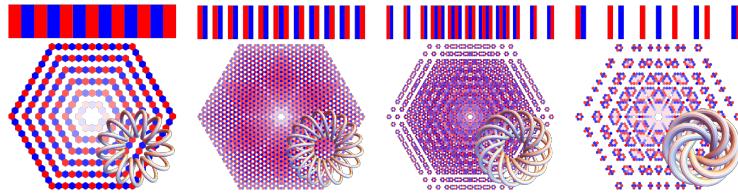
So θ detects knot mutation and separates the Conway knot K11n34 from the Kinoshita-Terasaka knot K11n42!



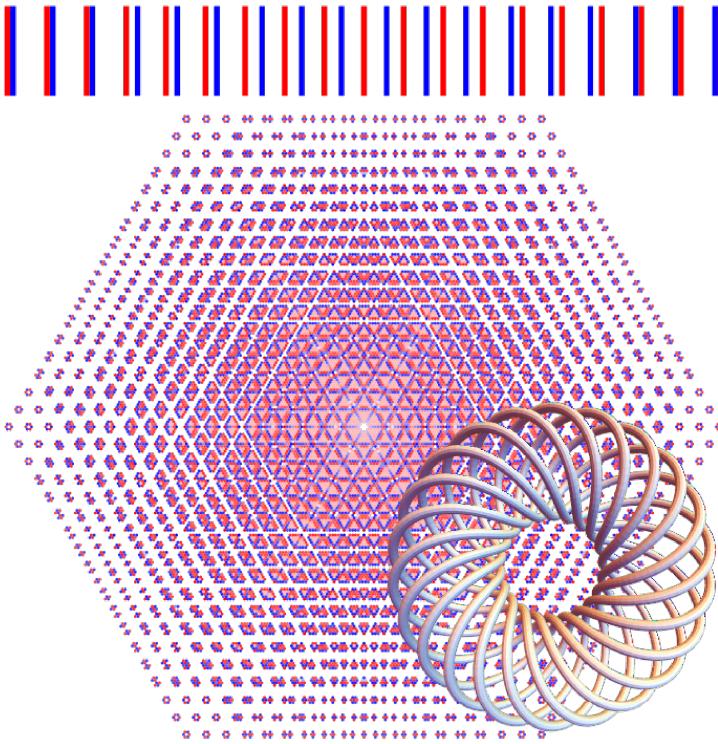
A faster program, in which the Feynman diagrams are “pre-computed” (see theta.nb at [ωεβ/ap](#)):

⊕ GraphicsRow[ImageCompose[

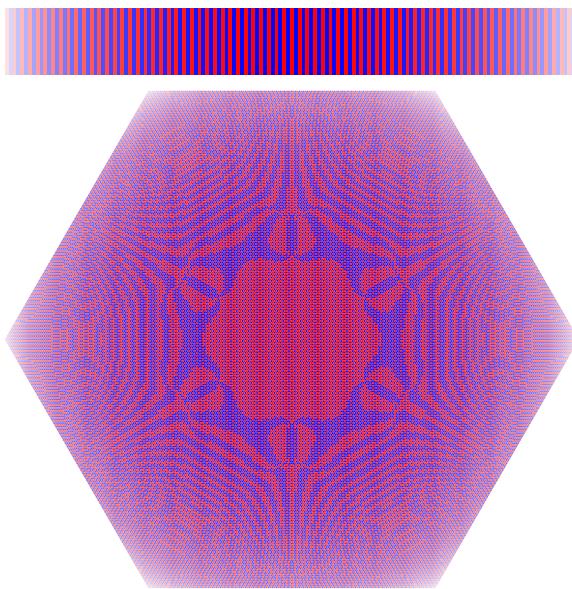
```
PolyPlot[θ[TorusKnot @@ #], ImageSize -> 480],
TubePlot[TorusKnot @@ #, ImageSize -> 240],
{Right, Bottom}, {Right, Bottom}
] & /@ {{13, 2}, {17, 3}, {13, 5}, {7, 6}}]
```



The torus knot $T_{22/7}$:



Next, a random 300 crossing knot from [DHOEBL] (more at [weβ/DK](#)):

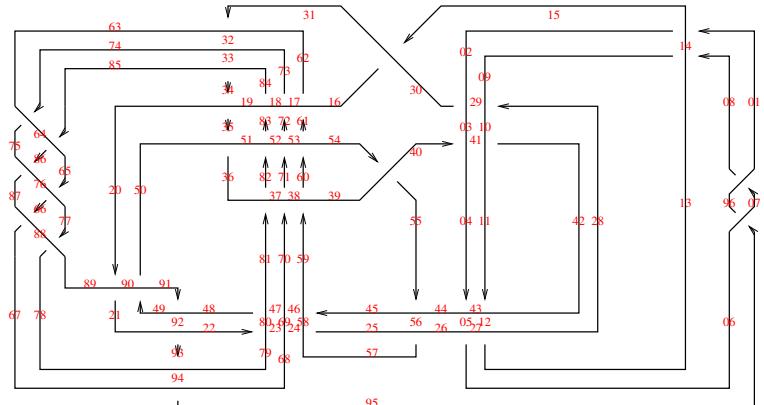


Unproven Fact. For any knot K , twice its genus $g(K)$ bounds the T_1 degree of θ : $\deg_{T_1} \theta(K) \leq 2g(K)$.

The 48-crossing Gompf-Scharlemann-Thompson GST_{48} knot [GST] is significant because it may be a counterexample to the slice-ribbon conjecture:

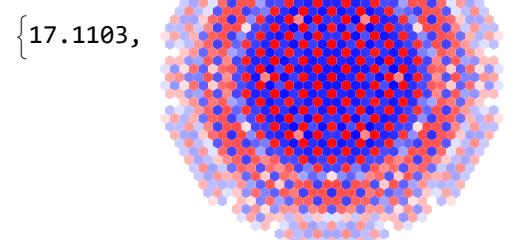


Gompf Scharlemann Thompson



⊕ $GST_{48} = EPD[X_{14,1}, \bar{X}_{2,29}, X_{3,40}, X_{43,4}, \bar{X}_{26,5}, X_{6,95}, X_{96,7}, X_{13,8}, \bar{X}_{9,28}, X_{10,41}, X_{42,11}, \bar{X}_{27,12}, X_{30,15}, \bar{X}_{16,61}, \bar{X}_{17,72}, \bar{X}_{18,83}, X_{19,34}, \bar{X}_{89,20}, \bar{X}_{21,92}, \bar{X}_{79,22}, \bar{X}_{68,23}, \bar{X}_{57,24}, \bar{X}_{25,56}, X_{62,31}, X_{73,32}, X_{84,33}, \bar{X}_{50,35}, X_{36,81}, X_{37,70}, X_{38,59}, \bar{X}_{39,54}, X_{44,55}, X_{58,45}, X_{69,46}, X_{80,47}, X_{48,91}, X_{90,49}, X_{51,82}, X_{52,71}, X_{53,60}, \bar{X}_{63,74}, \bar{X}_{64,85}, \bar{X}_{76,65}, \bar{X}_{87,66}, \bar{X}_{67,94}, \bar{X}_{75,86}, \bar{X}_{88,77}, \bar{X}_{78,93}]$;

AbsoluteTiming[PolyPlot[θ_{48} = θ@GST_{48}, ImageSize -> Small]]



⊕ {Exponent[θ_{48}[1], T], Floor[Exponent[θ_{48}[2], T_2] / 2]}

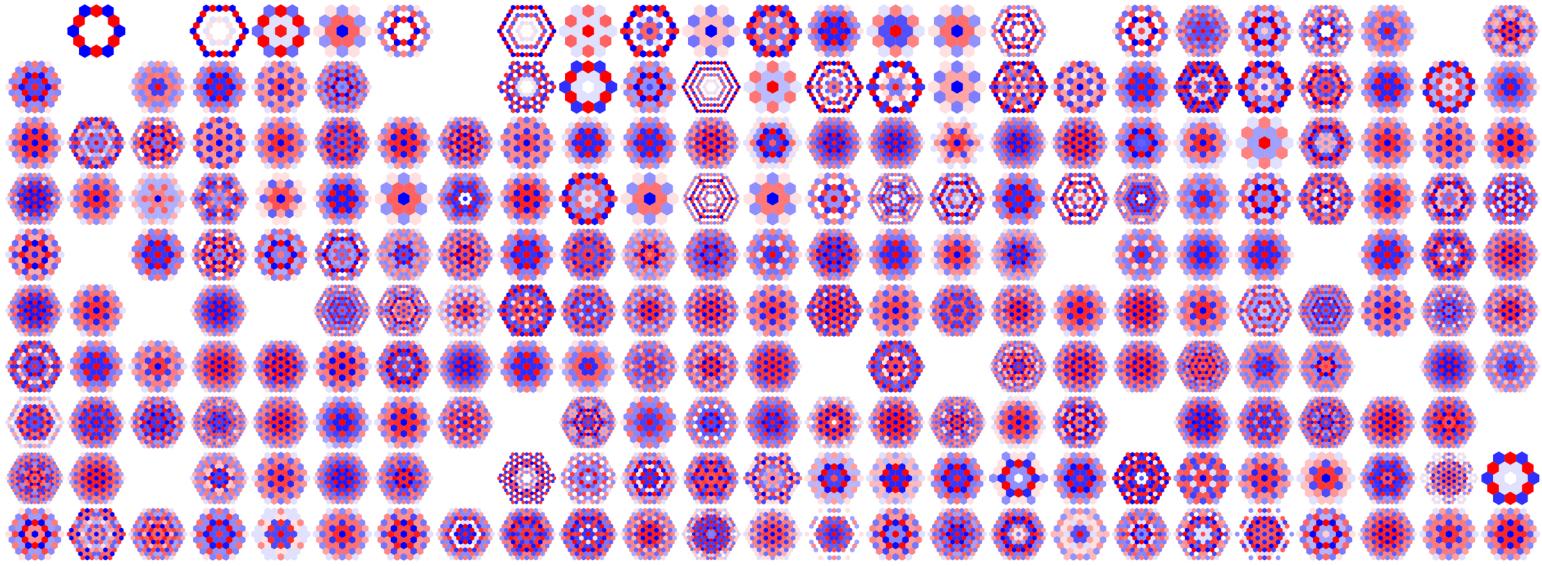
⊖ {8, 10}

So Θ knows things about GST_{48} that Δ doesn't!



Prior Art. θ is probably equal to the “2-loop polynomial” studied by Ohtsuki [Oh2] (at greater difficulty, with harder computations), continuing B-N, Garoufalidis, Rozansky, Kricker, and Schaveling [BNG, GR, R1, R2, R3, Kr, Sch]. θ is related, but probably not equivalent, to the invariant studied by Garoufalidis–Kashaev [GK].

The Rolfsen Table of Knots.



Where is it coming from? The most honest answer is “we don’t know” (and that’s good!). The second most, “undetermined coefficients for an ansatz that made sense”. The ansatz comes from the following principles / earlier work:

Morphisms have generating functions. Indeed, there is an isomorphism

$$\mathcal{G}: \text{Hom}(\mathbb{Q}[x_i], \mathbb{Q}[y_j]) \rightarrow \mathbb{Q}[y_j][\xi_i],$$

and by PBW, many relevant spaces are polynomial rings, though only as vector spaces.

Composition is integration. Indeed, if $f \in \text{Hom}(\mathbb{Q}[x_i], \mathbb{Q}[y_j])$ and $g \in \text{Hom}(\mathbb{Q}[y_j], \mathbb{Q}[z_k])$, then

$$\mathcal{G}(g \circ f) = \int e^{-y\eta} f g d\eta$$

Use universal invariants. These take values in a universal enveloping algebra (perhaps quantized), and thus they are expressible as long compositions of generating functions. See [La, Oh1].

“Solvable approximation” \leadsto perturbed Gaussians. Let \mathfrak{g} be a semisimple Lie algebra, let \mathfrak{h} be its Cartan subalgebra, and let \mathfrak{b}^u and \mathfrak{b}^l be its upper and lower Borel subalgebras. Then \mathfrak{b}^u has a bracket β , and as the dual of \mathfrak{b}^l it also has a cobracket δ , and in fact, $\mathfrak{g} \oplus \mathfrak{h} \cong \text{Double}(\mathfrak{b}^u, \beta, \delta)$. Let $\mathfrak{g}_\epsilon^+ := \text{Double}(\mathfrak{b}^u, \beta, \epsilon\delta)$ (mod ϵ^{d+1} it is solvable for any d). Then by [BV3, BN1] (in the case of $\mathfrak{g} = sl_2$) all the interesting tensors of $\mathcal{U}(\mathfrak{g}_\epsilon^+)$ (quantized or not) are perturbed Gaussian with perturbation parameter ϵ with understood bounds on the degrees of the perturbations.

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[BN2] D. Bar-Natan, *Algebraic Knot Theory*, talk given in Sydney, September 2019. Handout and video at [oeβ/AKT](#).

[BN3] D. Bar-Natan, *Cars, Interchanges, Traffic Counters, and some Pretty Darned Good Knot Invariants*, talk given in “Using Quantum Invariants to do Interesting Topology”, Oaxaca, Mexico, October 2022. Handout and video at [oeβ/Cars](#).

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[BV1] D. Bar-Natan and R. van der Veen, *A Polynomial Time Knot Polynomial*, Proc. Amer. Math. Soc. **147** (2019) 377–397, [arXiv:1708.04853](#).

[BV2] D. Bar-Natan and R. van der Veen, *A Perturbed-Alexander Invariant*, Quantum Topology **15** (2024) 449–472, [oeβ/APAI](#).

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[La] R. J. Lawrence, *Universal Link Invariants using Quantum Groups*, Proc. XVII Int. Conf. on Diff. Geom. Methods in Theor. Phys., Chester, England, August 1988. World Scientific (1989) 55–63.

[LV] D. López Neumann and R. van der Veen, *Genus Bounds from Unrolled Quantum Groups at Roots of Unity*, [arXiv:2312.02070](#).

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Thanks for bearing with me!