



Cars, Interchanges, Traffic Counters, and some Pretty Darned Good Knot Invariants

More at ωεβ/APAI

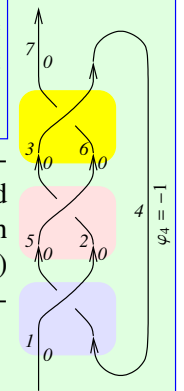
Abstract. Reporting on joint work with Roland van der Veen, I'll tell you some stories about ρ_1 , an easy to define, strong, fast to compute, homomorphic, and well-connected knot invariant. ρ_1 was first studied by Rozansky and Overbay [Ro1, Ro2, Ro3, Ov], it has far-reaching generalizations, it is dominated by the coloured Jones polynomial, and I wish I understood it. **Common misconception.** "Dominated" \Rightarrow "lesser".



Formulas stay; interpretations change with time.



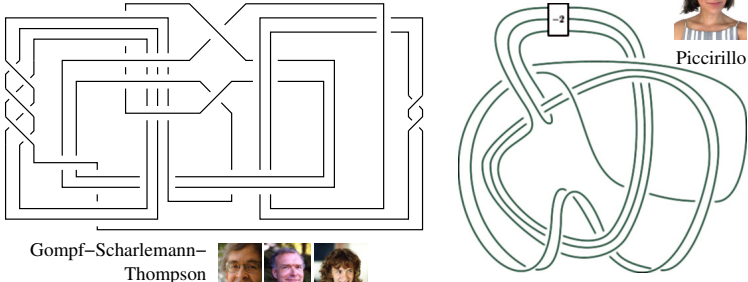
Dror Bar-Natan: Talks: CMS-2210: Simple, Concise, Powerful, and Not Understood. Abstract. I will give a simple and concise description of a strong and fast to compute knot invariant which seems to actually see some topology. Your homework will be to explain why it makes sense for a topologist to invert a presentation matrix for the Alexander module of a knot and consider quadratic combinations of the entries of said inverse. Thanks for inviting me to speak at this CMS meeting!



We seek strong, fast, homomorphic knot and tangle invariants.

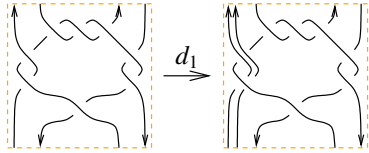
Strong. Having a small "kernel".

Fast. Computable even for large knots (best: poly time).

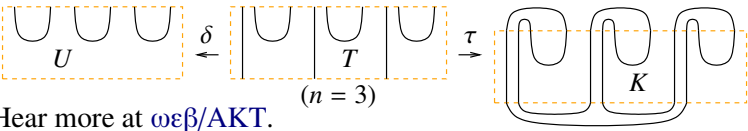


Gompf-Scharlemann-Thompson

Homomorphic. Extends to tangles and behaves under tangle operations; especially gluings and doublings:



Why care for "Homomorphic"? **Theorem.** A knot K is ribbon iff there exists a $2n$ -component tangle T with skeleton as below such that $\tau(T) = K$ and where $\delta(T) = U$ is the *untangle*:

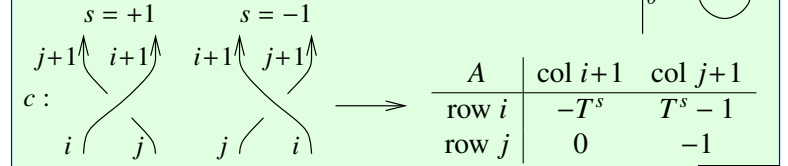


Hear more at ωεβ/AKT.

Acknowledgement. This work was partially supported by NSERC grant RGPIN-2018-04350 and by the Chu Family Foundation (NYC).

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Formulas. Draw an n -crossing knot K as on the right: all crossings face up, and the edges are marked with a running index $k \in \{1, \dots, 2n + 1\}$ and with rotation numbers φ_k . Let A be the $(2n + 1) \times (2n + 1)$ matrix constructed by starting with the identity matrix I , and adding a 2×2 block for each crossing:



Let $G = (g_{\alpha\beta}) = A^{-1}$. For the trefoil example, it is:

A	col $i+1$	col $j+1$
row i	$-T^s$	$T^s - 1$
row j	0	-1

Burau

Alexander

Fox

Wirtinger

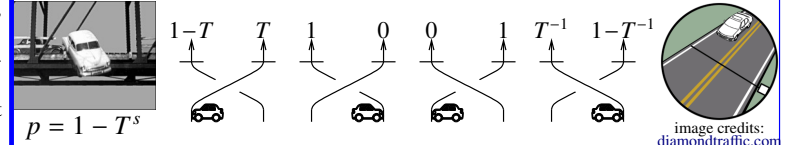
Blanchfield

"The Green Function"

Note. The Alexander polynomial Δ is given by $\Delta = T^{-(\varphi-w)/2} \det(A)$, with $\varphi = \sum_k \varphi_k$, $w = \sum_c s$.

Classical Topologists: This is boring. Yawn.
Formulas, continued. Finally, set $R_1(c) := s(g_{ji}(g_{j+1,j} + g_{j,j+1} - g_{ij}) - g_{ii}(g_{j,j+1} - 1) - 1/2)$
 $\rho_1 := \Delta^2 \left(\sum_c R_1(c) - \sum_k \varphi_k (g_{kk} - 1/2) \right)$.
 In our example $\rho_1 = -T^2 + 2T - 2 + 2T^{-1} - T^{-2}$.
Theorem. ρ_1 is a knot invariant. **Proof:** later. not today
Classical Topologists: Whiskey Tango Foxtrot?

Cars, Interchanges, and Traffic Counters. Cars always drive forward. When a car crosses over a bridge it goes through with (algebraic) probability $T^s \sim 1$, but falls off with probability $1 - T^s \sim 0^*$. See also [Jo, LTW].



* In algebra $x \sim 0$ if for every y in the ideal generated by x , $1 - y$ is invertible.

Preliminaries

This is Rho.nb of <http://drorbn.net/oa22/ap>.

Once [`<< KnotTheory``; `<< Rot.m`];

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of February 2, 2020, 10:53:45.2097.

Read more at <http://katlas.org/wiki/KnotTheory>.

Loading Rot.m from <http://drorbn.net/la22/ap>

to compute rotation numbers.

The Program

```

R1[s_, i_, j_] :=
  S (g_{ji} (g_{j^+, j} + g_{j, j^+} - g_{ij}) - g_{ii} (g_{j, j^+} - 1) - 1/2);
Z[K_] := Module[{Cs, phi, n, A, s, i, j, k, Delta, G, rho1},
  {Cs, phi} = Rot[K]; n = Length[Cs];
  A = IdentityMatrix[2 n + 1];
  Cases[Cs, {s_, i_, j_} ->
    (A[[{i, j}, {i + 1, j + 1}] + = ( -T^S T^S - 1 ) )];
  Delta = T^(-Total[phi] - Total[Cs[[All, 1]]) / 2 Det[A];
  G = Inverse[A];
  rho1 = Sum_{k=1}^n R1 @@ Cs[[k]] - Sum_{k=1}^{2n} phi[[k]] (g_{kk} - 1/2);
  Factor@
    {Delta, Delta^2 rho1 /. alpha_+ -> alpha + 1 /. g_{alpha, beta} -> G[[alpha, beta]]};

```

The First Few Knots

```

TableForm[Table[Join[{K[[1]]_k[[2]]}, Z[K]],
  {K, AllKnots[{3, 6]}], TableAlignments -> Center}]

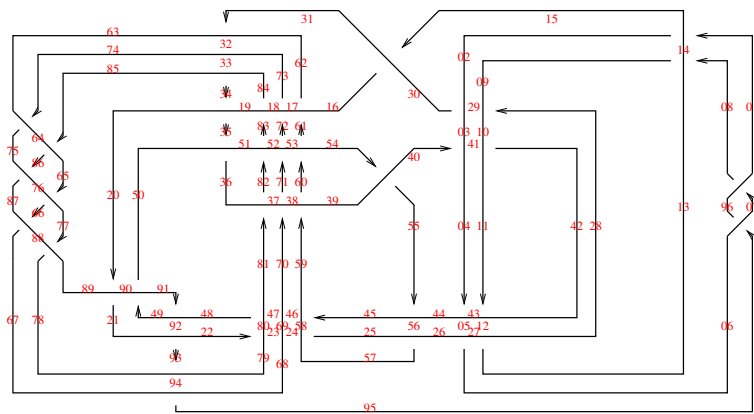
```

3_1	$\frac{1-T+T^2}{T}$	$\frac{(-1+T)^2(1+T^2)}{T^2}$
4_1	$-\frac{1-3T+T^2}{T}$	0
5_1	$\frac{1-T+T^2-T^3+T^4}{T^2}$	$\frac{(-1+T)^2(1+T^2)(2+T^2+2T^4)}{T^4}$
5_2	$\frac{2-3T+2T^2}{T}$	$\frac{(-1+T)^2(5-4T+5T^2)}{T^2}$
6_1	$-\frac{(-2+T)(-1+2T)}{T}$	$\frac{(-1+T)^2(1-4T+T^2)}{T^2}$
6_2	$-\frac{1-3T+3T^2-3T^3+T^4}{T^2}$	$\frac{(-1+T)^2(1-4T+4T^2-4T^3+4T^4-4T^5+T^6)}{T^4}$
6_3	$\frac{1-3T+5T^2-3T^3+T^4}{T^2}$	0



$$p = 1 - T^S$$

Fast!



Timing@

```

Z[GST48 = EPD[X_{14,1}, X_{2,29}, X_{3,40}, X_{43,4}, X_{26,5}, X_{6,95},
  X_{96,7}, X_{13,8}, X_{9,28}, X_{10,41}, X_{42,11}, X_{27,12}, X_{30,15},
  X_{16,61}, X_{17,72}, X_{18,83}, X_{19,34}, X_{89,20}, X_{21,92},
  X_{79,22}, X_{68,23}, X_{57,24}, X_{25,56}, X_{62,31}, X_{73,32},
  X_{84,33}, X_{50,35}, X_{36,81}, X_{37,70}, X_{38,59}, X_{39,54}, X_{44,55},
  X_{58,45}, X_{69,46}, X_{80,47}, X_{48,91}, X_{90,49}, X_{51,82}, X_{52,71},
  X_{53,60}, X_{63,74}, X_{64,85}, X_{76,65}, X_{87,66}, X_{67,94},
  X_{75,86}, X_{88,77}, X_{78,93}]]

```

$$\{170.313, \left\{ -\frac{1}{T^8} (-1 + 2T - T^2 - T^3 + 2T^4 - T^5 + T^8) \right.$$

$$\left. (-1 + T^3 - 2T^4 + T^5 + T^6 - 2T^7 + T^8) \right\}, \frac{1}{T^{16}}$$

$$(-1 + T)^2 (5 - 18T + 33T^2 - 32T^3 + 2T^4 + 42T^5 - 62T^6 - 8T^7 + 166T^8 - 242T^9 + 108T^{10} + 132T^{11} - 226T^{12} + 148T^{13} - 11T^{14} - 36T^{15} - 11T^{16} + 148T^{17} - 226T^{18} + 132T^{19} + 108T^{20} - 242T^{21} + 166T^{22} - 8T^{23} - 62T^{24} + 42T^{25} + 2T^{26} - 32T^{27} + 33T^{28} - 18T^{29} + 5T^{30}) \}$$

Strong!

```
{NumberOfKnots[{3, 12}],
```

```
Length@
```

```
Union@Table[Z[K], {K, AllKnots[{3, 12]}],
```

```
Length@
```

```
Union@Table[{HOMFLYPT[K], Kh[K]},
```

```
{K, AllKnots[{3, 12]}]}}
```

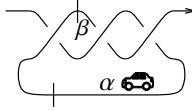
{2977, 2882, 2785}

So the pair (Δ, ρ_1) attains 2,882 distinct values on the 2,977 prime knots with up to 12 crossings (a deficit of 95), whereas the pair (HOMFLYPT, Khovanov Homology) attains only 2,785 distinct values on the same knots (a deficit of 192).

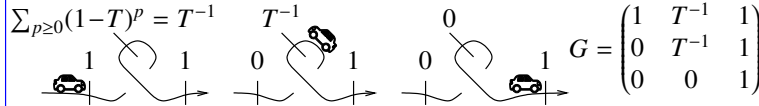


Hoste Ocneanu Millett Freyd Lickorish Yetter Przytycki Traczyk Khovanov

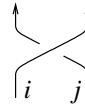
Theorem. The Green function $g_{\alpha\beta}$ is the reading of a traffic counter at β , if car traffic is injected at α (if $\alpha = \beta$, the counter is *after* the injection point).



Example.



Proof. Near a crossing c with sign s , incoming upper edge i and incoming lower edge j , both sides satisfy the g -rules:



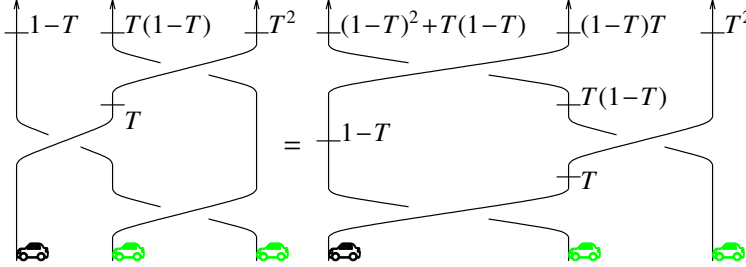
$$g_{i\beta} = \delta_{i\beta} + T^s g_{i+1,\beta} + (1 - T^s) g_{j+1,\beta}, \quad g_{j\beta} = \delta_{j\beta} + g_{j+1,\beta},$$

and always, $g_{\alpha,2n+1} = 1$: use common sense and $AG = I (= GA)$.

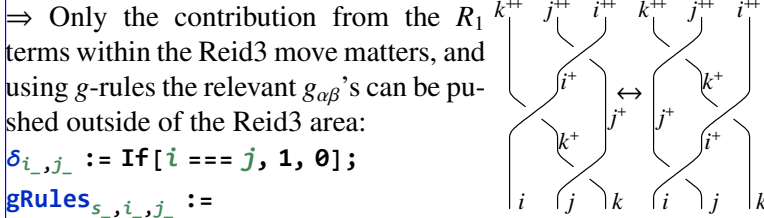
Bonus. Near c , both sides satisfy the further g -rules:

$$g_{\alpha i} = T^{-s}(g_{\alpha,i+1} - \delta_{\alpha,i+1}), \quad g_{\alpha j} = g_{\alpha,j+1} - (1 - T^s)g_{\alpha i} - \delta_{\alpha,j+1}.$$

Invariance of ρ_1 . We start with the hardest, Reidemeister 3:



⇒ Overall traffic patterns are unaffected by Reid3!
 ⇒ Green's $g_{\alpha\beta}$ is unchanged by Reid3, provided the cars injection site α and the traffic counters β are away.
 ⇒ Only the contribution from the R_1 terms within the Reid3 move matters, and using g -rules the relevant $g_{\alpha\beta}$'s can be pushed outside of the Reid3 area:



$$\delta_{i,j} := \text{If}[i == j, 1, 0];$$

$$\text{gRules}_{s,i,j} := \{ g_{i\beta} \mapsto \delta_{i\beta} + T^s g_{i+1,\beta} + (1 - T^s) g_{j+1,\beta}, g_{j\beta} \mapsto \delta_{j\beta} + g_{j+1,\beta}, g_{\alpha,i} \mapsto T^{-s}(g_{\alpha,i+1} - \delta_{\alpha,i+1}), g_{\alpha,j} \mapsto g_{\alpha,j+1} - (1 - T^s)g_{\alpha i} - \delta_{\alpha,j+1} \}$$

$$\text{lhs} = R_1[1, j, k] + R_1[1, i, k^+] + R_1[1, i^+, j^+] // .$$

$$\text{gRules}_{1,j,k} \cup \text{gRules}_{1,i,k^+} \cup \text{gRules}_{1,i^+,j^+};$$

$$\text{rhs} = R_1[1, i, j] + R_1[1, i^+, k] + R_1[1, j^+, k^+] // .$$

$$\text{gRules}_{1,i,j} \cup \text{gRules}_{1,i^+,k} \cup \text{gRules}_{1,j^+,k^+};$$

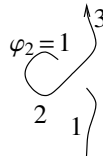
Simplify[lhs == rhs]

True

Next comes Reid1, where we use results from an earlier example:

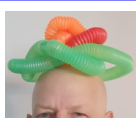
$$R_1[1, 2, 1] - 1 (g_{22} - 1/2) /. g_{\alpha,\beta} \mapsto \begin{pmatrix} 1 & T^{-1} & 1 \\ 0 & T^{-1} & 1 \\ 0 & 0 & 1 \end{pmatrix} [[\alpha, \beta]]$$

$$\frac{1}{T^2} - \frac{1}{T} - \frac{-1 + \frac{1}{T}}{T} = 0$$



Invariance under the other moves is proven similarly.

Wearing my Topology hat the formula for R_1 , and even the idea to look for R_1 , remain a complete mystery to me.



Wearing my Quantum Algebra hat, I spy a Heisenberg algebra $\mathbb{H} = A\langle p, x \rangle / ([p, x] = 1)$:

$$\text{cars} \leftrightarrow p \quad \text{traffic counters} \leftrightarrow x$$



Where did it come from? Consider $g_\epsilon := sl_{2+}^\epsilon := L\langle y, b, a, x \rangle$ with relations

$$[b, x] = \epsilon x, \quad [b, y] = -\epsilon y, \quad [b, a] = 0,$$

$$[a, x] = x, \quad [a, y] = -y, \quad [x, y] = b + \epsilon a.$$

At invertible ϵ , it is isomorphic to sl_2 plus a central factor, and it can be quantized à la Drinfel'd [Dr] much like sl_2 to get an algebra $QU = A\langle y, b, a, x \rangle$ subject to (with $q = e^{\hbar\epsilon}$):

$$[b, a] = 0, \quad [b, x] = \epsilon x, \quad [b, y] = -\epsilon y,$$

$$[a, x] = x, \quad [a, y] = -y, \quad xy - qyx = \frac{1 - e^{-\hbar(b+\epsilon a)}}{\hbar}.$$

Now QU has an R -matrix solving Yang-Baxter (meaning Reid3),

$$R = \sum_{m,n \geq 0} \frac{y^m b^m \otimes (\hbar a)^m (\hbar x)^n}{m! [n]_q!}, \quad ([n]_q! \text{ is a "quantum factorial"})$$

and so it has an associated "universal quantum invariant" à la Lawrence and Ohtsuki [La, Oh], $Z_\epsilon(K) \in QU$.

Now $QU \cong \mathcal{U}(g_\epsilon)$ (only as algebras!) and $\mathcal{U}(g_\epsilon)$ represents into \mathbb{H} via

$$y \rightarrow -tp - \epsilon \cdot xp^2, \quad b \rightarrow t + \epsilon \cdot xp, \quad a \rightarrow xp, \quad x \rightarrow x,$$

(abstractly, g_ϵ acts on its Verma module

$$\mathcal{U}(g_\epsilon) / (\mathcal{U}(g_\epsilon)\langle y, a, b - \epsilon a - t \rangle) \cong \mathbb{Q}[x]$$

by differential operators, namely via \mathbb{H}), so R can be pushed to $\mathcal{R} \in \mathbb{H} \otimes \mathbb{H}$.

Everything still makes sense at $\epsilon = 0$ and can be expanded near $\epsilon = 0$ resulting with $\mathcal{R} = \mathcal{R}_0(1 + \epsilon \mathcal{R}_1 + \dots)$, with $\mathcal{R}_0 = e^{t(xp \otimes 1 - x \otimes p)}$ and \mathcal{R}_1 a quartic polynomial in p and x . So p 's and x 's get created along K and need to be pushed around to a standard location ("normal ordering"). This is done using

$$(p \otimes 1)\mathcal{R}_0 = \mathcal{R}_0(T(p \otimes 1) + (1 - T)(1 \otimes p)),$$

$$(1 \otimes p)\mathcal{R}_0 = \mathcal{R}_0(1 \otimes p),$$

and when the dust settles, we get our formulas for ρ_1 . But QU is a quasi-triangular Hopf algebra, and hence ρ_1 is **homomorphic**. Read more at [BV1, BV2] and hear more at $\omega\epsilon\beta/\text{SolvApp}$, $\omega\epsilon\beta/\text{Dogma}$, $\omega\epsilon\beta/\text{DoPeGDO}$, $\omega\epsilon\beta/\text{FDA}$, $\omega\epsilon\beta/\text{AQDW}$.

Also, we can (and know how to) look at higher powers of ϵ and we can (and more or less know how to) replace sl_2 by arbitrary semi-simple Lie algebra (e.g., [Sch]). So ρ_1 is **not alone!**



Schaveling

These constructions are very similar to Rozansky-Overbay [Ro1, Ro2, Ro3, Ov] and hence to the "loop expansion" of the Kontsevich integral and the coloured Jones polynomial.

If this all reads like **insanity** to you, it should (and you haven't seen half of it). Simple things should have simple explanations.

Hence, **Homework**. Explain ρ_1 with no reference to quantum voodoo and find it a topology home (large enough to house generalizations!). Make explicit the homomorphic properties of ρ_1 . Use them to do topology!

P.S. As a friend of Δ , ρ_1 gives a genus bound, sometimes better than Δ 's. How much further does this friendship extend?

A Small-Print Page on $\rho_d, d > 1$.

Definition. $\langle f(z_i), h(\zeta_i) \rangle_{\{z_i\}} := f(\partial_{\zeta_i} h)|_{\zeta_i=0}$, so $\langle p^2 x^2, \mathbb{Q}^{g\pi\xi} \rangle = 2g^2$.

Baby Theorem. There exist (non unique) power series $r^\pm(p_1, p_2, x_1, x_2) = \sum_d \epsilon^d r_d^\pm(p_1, p_2, x_1, x_2) \in \mathbb{Q}[T^{\pm 1}, p_1, p_2, x_1, x_2][[\epsilon]]$ with $\deg r_d^\pm \leq 2d + 2$ ("docile") such that the power series $Z^b = \sum \rho_d \epsilon^d :=$

$$\left\langle \exp\left(\sum_c r^s(p_i, p_j, x_i, x_j)\right), \exp\left(\sum_{\alpha\beta} g_{\alpha\beta} \pi_\alpha \xi_\beta\right) \right\rangle_{\{p_\alpha, x_\beta\}}$$

is a bnot invariant. Beyond the once-and-for-all computation of $g_{\alpha\beta}$ (a matrix inversion), Z^b is computable in $O(n^d)$ operations in the ring $\mathbb{Q}[T^{\pm 1}]$.

(Bnots are knot diagrams modulo the braid-like Reidemeister moves, but not the cyclic ones).

Theorem. There also exist docile power series $\gamma^\varphi(\bar{p}, \bar{x}) = \sum_d \epsilon^d \gamma_d^\varphi \in \mathbb{Q}[T^{\pm 1}, \bar{p}, \bar{x}][[\epsilon]]$ such that the power series $Z = \sum \rho_d \epsilon^d :=$

$$\left\langle \exp\left(\sum_c r^s(p_i, p_j, x_i, x_j) + \sum_k \gamma^{\varphi k}(\bar{p}_k, \bar{x}_k)\right), \exp\left(\sum_{\alpha\beta} g_{\alpha\beta}(\pi_\alpha + \bar{\pi}_\alpha)(\xi_\beta + \bar{\xi}_\beta) + \sum_\alpha \pi_\alpha \bar{\xi}_\alpha\right) \right\rangle_{\{p_\alpha, \bar{p}_\alpha, x_\beta, \bar{x}_\beta\}}$$

is a knot invariant, as easily computable as Z^b .

Implementation. Data, then program (with output using the Conway variable $z = \sqrt{T} - 1 / \sqrt{T}$), and then a demo. See `Rho.nb` of `omegaBeta/alpha`.

`V@r_{1,\varphi}[k_] := \varphi (1/2 - \bar{p}_k \bar{x}_k)` ; `V@r_{2,\varphi}[k_] := -\varphi^2 \bar{p}_k \bar{x}_k / 2` ;

`V@r_{3,\varphi}[k_] := -\varphi^3 \bar{p}_k \bar{x}_k / 6`

`V@r_{1,s}[i_, j_] := s (-1 + 2 p_i x_i - 2 p_j x_j + (-1 + T^s) p_i p_j x_i^2 + (1 - T^s) p_j^2 x_i^2 - 2 p_i p_j x_i x_j + 2 p_j^2 x_i x_j) / 2`

`V@r_{2,1}[i_, j_] := (-6 p_i x_i + 6 p_j x_j - 3 (-1 + 3 T) p_i p_j x_i^2 + 3 (-1 + 3 T) p_j^2 x_i^2 + 4 (-1 + T) p_i^2 p_j x_i^2 - 2 (-1 + T) (5 + T) p_i p_j^2 x_i^2 + 2 (-1 + T) (3 + T) p_j^3 x_i^2 + 18 p_i p_j x_i x_j - 18 p_j^2 x_i x_j - 6 p_i^2 p_j x_i x_j + 6 (2 + T) p_i p_j^2 x_i x_j - 6 (1 + T) p_j^3 x_i x_j - 6 p_i p_j^2 x_i x_j^2 + 6 p_j^3 x_i x_j^2) / 12`

`V@r_{2,-1}[i_, j_] := (-6 T^2 p_i x_i + 6 T^2 p_j x_j + 3 (-3 + T) T p_i p_j x_i^2 - 3 (-3 + T) T p_j^2 x_i^2 - 4 (-1 + T) T p_i^2 p_j x_i^2 + 2 (-1 + T) (1 + 5 T) p_i p_j^2 x_i^2 - 2 (-1 + T) (1 + 3 T) p_j^3 x_i^2 + 18 T^2 p_i p_j x_i x_j - 18 T^2 p_j^2 x_i x_j - 6 T^2 p_i^2 p_j x_i x_j + 6 T (1 + 2 T) p_i p_j^2 x_i x_j - 6 T (1 + T) p_j^3 x_i x_j - 6 T^2 p_i p_j^2 x_i x_j^2 + 6 T^2 p_j^3 x_i x_j^2) / (12 T^2)`

`Z_2[GST48] (* takes a few minutes *)`

$$\{1 - 4 z^2 - 61 z^4 - 207 z^6 - 296 z^8 - 210 z^{10} - 77 z^{12} - 14 z^{14} - z^{16}, 1 + (38 z^2 + 255 z^4 + 1696 z^6 + 16281 z^8 + 86952 z^{10} + 259994 z^{12} + 487372 z^{14} + 615066 z^{16} + 543148 z^{18} + 341714 z^{20} + 153722 z^{22} + 48983 z^{24} + 10776 z^{26} + 1554 z^{28} + 132 z^{30} + 5 z^{32}) e + (-8 - 484 z^2 + 9709 z^4 + 165952 z^6 + 1590491 z^8 + 16256508 z^{10} + 115341797 z^{12} + 432685748 z^{14} + 395838354 z^{16} - 4017557792 z^{18} - 23300064167 z^{20} - 70082264972 z^{22} - 142572271191 z^{24} - 209475503700 z^{26} - 221616295209 z^{28} - 151502648428 z^{30} - 23700199243 z^{32} + 99462146328 z^{34} + 164920463074 z^{36} + 162550825432 z^{38} + 119164552296 z^{40} + 69153062608 z^{42} + 32547596611 z^{44} + 12541195448 z^{46} + 3961384155 z^{48} + 1021219696 z^{50} + 212773106 z^{52} + 35264208 z^{54} + 4537548 z^{56} + 436600 z^{58} + 29536 z^{60} + 1252 z^{62} + 25 z^{64}) e^2\}$$

`TableForm[Table[Join[{K[[1]][K[[2]]], Z_3[K]}, {K, AllKnots[{3, 6]}}], TableAlignments -> Center] (* takes a few minutes *)`

3_1	$1 + z^2$	$1 + (2z^2 + z^4) e + (2 - 4z^2 + 3z^4 + 4z^6 + z^8) e^2 + (-12 + 74z^2 - 27z^4 - 20z^6 + 8z^8 + 6z^{10} + z^{12}) e^3$
4_1	$1 - z^2$	$1 + (-2 - 2z^4) e^2$
5_1	$1 + 3z^2 + z^4$	$1 + (10z^2 + 21z^4 + 12z^6 + 2z^8) e + (6 - 28z^2 + 33z^4 + 364z^6 + 655z^8 + 536z^{10} + 227z^{12} + 48z^{14} + 4z^{16}) e^2 + (-60 - 970z^2 + 645z^4 - 3380z^6 - 3280z^8 - 7470z^{10} - 19475z^{12} + 20536z^{14} + 12564z^{16} + 4774z^{18} + 1109z^{20} + 144z^{22} + 8z^{24}) e^3$
5_2	$1 + 2z^2$	$1 + (6z^2 + 5z^4) e + (4 - 20z^2 + 43z^4 + 64z^6 + 26z^8) e^2 + (-36 + 498z^2 - 883z^4 + 100z^6 + 816z^8 + 556z^{10} + 146z^{12}) e^3$
6_1	$1 - 2z^2$	$1 + (-2z^2 + z^4) e + (-4 + 4z^2 + 25z^4 - 8z^6 + 2z^8) e^2 + (12 + 154z^2 - 223z^4 - 608z^6 + 100z^8 - 52z^{10} + 10z^{12}) e^3$
6_2	$1 - z^2 - z^4$	$1 + (-2z^2 - 3z^4 + 2z^6 + z^8) e + (-2 - 4z^2 + 29z^4 + 28z^6 + 42z^8 - 8z^{10} - 2z^{12} + 4z^{14} + z^{16}) e^2 + (12 + 166z^2 + 155z^4 - 194z^6 - 2453z^8 - 1622z^{10} - 1967z^{12} - 258z^{14} + 49z^{16} - 30z^{18} + z^{20} + 6z^{22} + 2z^{24}) e^3$
6_3	$1 + z^2 + z^4$	$1 + (2 + 8z^2 - 16z^4 - 24z^6 - 16z^{10} - 2z^{12}) e^2$

`V@r_{3,1}[i_, j_] := (4 p_i x_i - 4 p_j x_j + 2 (5 + 7 T) p_i p_j x_i^2 - 2 (5 + 7 T) p_j^2 x_i^2 - 4 (-5 + 6 T) p_i^2 p_j x_i^2 + 4 (-16 + 17 T + 2 T^2) p_i p_j^2 x_i^2 - 4 (-11 + 11 T + 2 T^2) p_j^3 x_i^2 + 3 (-1 + T) p_i^3 p_j x_i^2 - 3 (-1 + T) (4 + 3 T) p_i^2 p_j^2 x_i^2 + (-1 + T) (13 + 22 T + T^2) p_i p_j^3 x_i^2 - (-1 + T) (4 + 13 T + T^2) p_j^4 x_i^2 - 28 p_i p_j x_i x_j + 28 p_j^2 x_i x_j + 36 p_i^2 p_j x_i x_j - 12 (9 + 2 T) p_i p_j^2 x_i x_j + 24 (3 + T) p_j^3 x_i x_j - 4 p_i^3 p_j x_i x_j + 28 T p_i^2 p_j^2 x_i x_j - 4 (-6 + 17 T + T^2) p_i p_j^3 x_i x_j + 4 (-5 + 10 T + T^2) p_j^4 x_i x_j + 24 p_i p_j^3 x_i x_j^2 - 24 p_i^2 p_j^2 x_i x_j^2 + 6 (10 + T) p_i p_j^3 x_i x_j^2 - 6 (6 + T) p_j^4 x_i x_j^2 - 4 p_i p_j^3 x_i x_j^2 + 4 p_j^4 x_i x_j^2) / 24`

`V@r_{3,-1}[i_, j_] := (-4 T^3 p_i x_i + 4 T^3 p_j x_j - 2 T^2 (7 + 5 T) p_i p_j x_i^2 + 2 T^2 (7 + 5 T) p_j^2 x_i^2 - 4 T^2 (-6 + 5 T) p_i^2 p_j x_i^2 + 4 T (-2 - 17 T + 16 T^2) p_i p_j^2 x_i^2 - 4 T (-2 - 11 T + 11 T^2) p_j^3 x_i^2 + 3 (-1 + T) T^2 p_i^3 p_j x_i^2 - 3 (-1 + T) T (3 + 4 T) p_i^2 p_j^2 x_i^2 + (-1 + T) (1 + 22 T + 13 T^2) p_i p_j^3 x_i^2 - (-1 + T) (1 + 13 T + 4 T^2) p_j^4 x_i^2 + 28 T^3 p_i p_j x_i x_j - 28 T^3 p_j^2 x_i x_j - 36 T^3 p_i^2 p_j x_i x_j + 12 T^2 (-2 + 9 T) p_i p_j^2 x_i x_j - 24 T^2 (1 + 3 T) p_j^3 x_i x_j + 4 T^3 p_i^3 p_j x_i x_j - 28 T^2 p_i^2 p_j^2 x_i x_j - 4 T (-1 - 17 T + 6 T^2) p_i p_j^3 x_i x_j + 4 T (-1 - 10 T + 5 T^2) p_j^4 x_i x_j - 24 T^3 p_i p_j^2 x_i x_j^2 + 24 T^3 p_j^3 x_i x_j^2 + 24 T^3 p_i^2 p_j x_i x_j^2 - 6 T^2 (1 + 10 T) p_i p_j^3 x_i x_j^2 + 6 T^2 (1 + 6 T) p_j^4 x_i x_j^2 + 4 T^3 p_i p_j^3 x_i x_j^2 - 4 T^3 p_j^4 x_i x_j^2) / (24 T^3)`

`{p*, x*, p-bar, x-bar} = {pi, xi, pi-bar, xi-bar}; (z-underline)_i := (z*)_i;`

`Zip_{i}[e_] := e;`

`Zip_{(z, z-underline)}[e_] :=`

`(Collect[e // Zip_{(zs)}, z] /. f_ . z^d_ -> {D[f, {z*, d}]} /. z* -> 0`

`gPair[f_s_, w_] :=`

`gPair[f_s, w] =`

`Collect[Zip_{Join@Table[{p_alpha, p-bar_alpha, x_alpha, x-bar_alpha}, {alpha, w}][`

`(Times @@ (V @ f_s))`

`Exp[Sum[g_{alpha, beta} (pi_alpha + pi-bar_alpha) (xi_beta + xi-bar_beta), {alpha, w}, {beta, w}] - Sum[xi-bar_alpha pi_alpha, {alpha, w}]]],`

`g_., Factor]`

`T2z[p_] := Module[{q = Expand[p], n, c},`

`If[q == 0, 0, c = Coefficient[q, T, n = Exponent[q, T]];`

`c z^n + T2z[q - c (T^{1/2} - T^{-1/2})^2 n]]];`

`Z_u[K_] := Module[{Cs, phi, n, A, s, i, j, k, Delta, G, d1, Z1, Z2, Z3},`

`{Cs, phi} = Rot[K]; n = Length[Cs]; A = IdentityMatrix[2 n + 1];`

`Cases[Cs, {s_., i_., j_} -> {A[[{i, j}], {i + 1, j + 1}]} + (-T^s T^s - 1)];`

`{Delta, G} = Factor@{T^{-Total[phi] - Total[Cs[[All, 1]]]} / 2 Det@A, Inverse@A};`

`Z1 =`

`Exp[Total[Cases[Cs, {s_., i_., j_} -> Sum[e^{d1} r_{d1, s}[i, j], {d1, d}]]] +`

`Sum[e^{d1} Y_{d1, phi[[k]]}[k], {k, 2 n}, {d1, d}] /. Y_u[e_] -> 0];`

`Z2 = Expand[F[{}], {}] x Normal@Series[Z1, {e, 0, d}] //.`

`F[f_s_, {e_s_...}] x (f : (r | Y)_{p_s_}[i_s_])^{p_s_} ->`

`F[Join[f_s, Table[f, p]], DeleteDuplicates@{e_s, i_s}];`

`Z3 = Expand[Z2 /. F[f_s_, e_s_] -> Expand[gPair[`

`Replace[f_s, Thread[e_s -> Range@Length@e_s], {2}], Length@e_s`

`] /. g_{alpha, beta} -> G[[e_s[[alpha], e_s[[beta]]]]];`

`Collect[{Delta, Z3 /. e^{p_s_} -> p! Delta^{p_s} e^{p_s}], e, T2z];`