

Equilibrium states:

- * something like a manifold
- * Equipped with "state functions": E, p, V, T, \dots
- * Empirical temperature scale

Digression: The ideal gas temp. scale:

$$pV = nRT$$

\uparrow # of moles \uparrow ideal gas constant

Experiments show that $\frac{pV}{n}$ approaches a limit (indep. of the type of gas) as the density goes to 0. That limit, divided by some constant R , is the empirical temp. scale.

Transitions between equilibrium states.

Example Gas in a cylinder with a movable piston.



- * Can constrain the volume V .
- * Can put a fixed weight on the piston to fix the pressure p .
- * Can put in a bath \rightarrow "temperature bath" to fix T .

Consider a transition $A \rightarrow B$ between equilibrium states.

The system does some amount W of work on the environment. The system also absorbs a certain amount Q of heat from the environment.

The first law of thermodynamics:

The First law of Thermodynamics:

$$E(B) - E(A) = Q - W$$

Remark E is a state function but Q & W are not.

Transition 1: Pull the piston so fast that the gas can't keep up. Then $W=0$. If everything is isolated, $E(B)=E(A)$.

Gay-Lussac did this on dilute gases, & found $T(B)=T(A)$.

\Rightarrow The energy of an ideal gas depends only on the temperature.

Transition 2: Expand very slowly, so that at all times the system is at equilibrium.

$$\text{force on piston} = F = \sum_{\substack{\uparrow \\ \text{area of piston}}} p(V) \quad \leftarrow \begin{array}{l} \text{pressure} \\ \text{at volume } V \end{array}$$

$$\text{So infinitesimal change in } W = \sum p(V) \cdot dy = p dV$$

$$\text{So } W = \int_{y_0}^{y_1} \underbrace{p(V) dV}_{\delta W}$$

So "work done" is the integral of a 1-form along a path.

This process is reversible.

In the same spirit there is a 1-form δQ s.t. for reversible transitions, $Q = \int \delta Q$.

So the first law, infinitesimally,

$$\delta E = \delta Q - \delta W \quad (\text{or } \delta Q = \delta E + p dV)$$

becomes
 $T dS = dE + p dV$

Imagine now pulling out the piston at an intermediate,

- "Finite" speed: - Complicated. Has to do w/ gas dynamics, sound waves, ...
- The system will not have a well-defined overall temperature.

So while the process is going on, the system is not in equilibrium.

- typically not reversible.

So the thermodynamic setup is:

- * A "manifold" of equilibrium states.
- * State functions: E, T, P, V, \dots
- * Two 1-forms: dW, dQ
- * A notion of "reversibility"

"Extensive": proportional to the amount of matter present. (E.g. V, E)

"Intensive": Independent of the amount of matter present. (E.g.: P, T)

"Specific" quantities: $v := V/N, e := E/N$ - these are intensive

The Second Law of Thermodynamics. A cyclic process is

$A_1 \rightarrow A_2 \rightarrow \dots \rightarrow A_m = A_1$ (seq. of processes)

The individual steps need not be reversible.

Kelvin's form: No thermodynamic system admits a cyclic process whose only effect is to extract a positive amount of heat from a bath and to convert it to work. [The other direction is]

possible!!

Suggestion Read about Carnot cycles in Fermi or in class notes

Consequence 1: Second law, reversible part:

There exists an empirical temperature scale T so that $\frac{\delta Q}{T}$ is exact; $\frac{\delta Q}{T} = dS$ where S is the "entropy".

This T is called "the absolute temp."

... at the end one shows that the absolute temp is equal to the ideal gas temp., up to a multiplicative scale.

In practice, δQ is always replaced with $T dS$

Consequence 2 Second law, irreversible part: If a thermally isolated system can make a transition from a state A to a state B , then $S(B) \geq S(A)$. If an amount Q of heat is absorbed at temp T , then

$$S(B) - S(A) \geq \frac{Q}{T}$$

regardless of work done.

Question 1. Is there a notion of a "thermodynamic manifold"?

2. Is there a notion of a "thermodynamic reality", a system of combining / breaking up thermodynamic manifolds.

δQ annihilates the tangent vector to γ , if γ is

Thermally isolated $\Rightarrow dS(i) = 0 \Rightarrow S(B) = S(A)$